

Change of basis

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We know that the eigenvectors of Hermitian operators may be chosen to form a complete, orthonormal set of basis states. Let

$$S = \{|a'\rangle \mid a' \in K\}$$

where K is some (finite or infinite) index set, be such a complete orthonormal basis, with

$$\hat{O}|a'\rangle = a'|a'\rangle$$

for some Hermitian operator, \hat{O} . There is a converse to this: if we have a complete, orthonormal set, then we can build operators:

$$\hat{O}' \equiv \sum_{a' \in K} \alpha(a') |a'\rangle \langle a'|$$

and the complex numbers $\alpha(a')$ will be its eigenvalues, since

$$\begin{aligned} \hat{O}'|a''\rangle &= \sum_{a' \in K} \alpha(a') |a'\rangle \langle a'| a''\rangle \\ &= \sum_{a' \in K} \alpha(a') |a'\rangle \delta_{a',a''} \\ &= \alpha(a'') |a''\rangle \end{aligned}$$

By choosing $\alpha(a')$ real for each a' , we insure that \hat{O}' is Hermitian. Therefore, we may find (many) Hermitian operators with a given complete orthonormal set of eigenstates.

If we have any quantum state, it may be expanded in term of any complete basis:

$$\begin{aligned} |\psi\rangle &= \hat{1}|\psi\rangle \\ &= \sum_{a' \in K} |a'\rangle \langle a'| \psi\rangle \end{aligned}$$

so the complex numbers $\langle a'| \psi\rangle$ represent the *state* $|\psi\rangle$ in the *basis* $|a'\rangle$.

If we wish to change the basis, consider the relationship between two bases, $|a'\rangle$ and $|b'\rangle$. Because these span the same vector space, there are equal numbers of indices, a' and b' and we put these in 1-1, but otherwise arbitrary, correspondence, $a_1 \leftrightarrow b_1, a_2 \leftrightarrow b_2, \dots$. In general, let $a' \leftrightarrow b'$. Now, since the basis ket $|b'\rangle$ is also a state, it may be expanded in terms of the $|a'\rangle$,

$$|b'\rangle = \sum_{a'' \in K} |a''\rangle \langle a''| b'\rangle$$

We may view this as a matrix \hat{U} , with components $\langle a''| b'\rangle$, acting on $|a''\rangle$ to give $|b'\rangle$. The matrix is the sum over corresponding pairs,

$$\hat{U} = \sum_m |b_m\rangle \langle a_m|$$

since then

$$\begin{aligned}\hat{U} |a_k\rangle &= \sum_m |b_m\rangle \langle a_m | a_k\rangle \\ &= \sum_m \delta_{mk} |b_m\rangle \\ &= |b_k\rangle\end{aligned}$$

Using this relation we see that this operator indeed has the right matrix components

$$\langle a_l | \hat{U} |a_k\rangle = \langle a_l | b_k\rangle$$

Notice that \hat{U} is unitary, since, with $\hat{U}^\dagger = \sum_m |a_m\rangle \langle b_m|$, we have

$$\begin{aligned}\hat{U}^\dagger \hat{U} &= \left(\sum_m |a_m\rangle \langle b_m| \right) \left(\sum_n |b_n\rangle \langle a_n| \right) \\ &= \sum_{m,n} |a_m\rangle \langle b_m | b_n\rangle \langle a_n| \\ &= \sum_{m,n} \delta_{mn} |a_m\rangle \langle a_n| \\ &= \sum_m |a_m\rangle \langle a_m| \\ &= \hat{1}\end{aligned}$$