

# Problems in the Schwarzschild geometry

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Start work on the problems below. See the notes on Schwarzschild geodesics to get started. Help is available.

## Summary of geodesics:

We have, for timelike curves:

$$\begin{aligned}\frac{dt}{d\tau} &= u_0^0 \left( \frac{1 - \frac{2M}{r_0}}{1 - \frac{2M}{r}} \right) \\ \frac{d\varphi}{d\tau} &= \frac{L}{r^2} \\ \frac{dr}{d\tau} &= \sqrt{2E + \frac{2M}{r} - \frac{L^2}{r^2} + \frac{2ML^2}{r^3}}\end{aligned}$$

where

$$\begin{aligned}L &= r_0^2 \left( \frac{d\varphi}{d\tau} \right)_0 \\ E &= -\frac{1}{2} \left[ 1 - \left( 1 - \frac{2M}{r_0} \right)^2 (u_0^0)^2 \right]\end{aligned}$$

## Problems:

1. Find the form of  $\frac{dr}{d\tau}$  for spacelike and null geodesics.
2. Find the proper time required for a particle to fall from radius  $r_0 > 2M$  to  $r = 2M$ . Evaluate the time numerically for a solar mass black hole, and a galactic black hole with mass  $10^7$  times the mass of the sun.
3. Find the proper time required for a particle to fall from radius  $r_0 = 2M$  to  $r = 0$ . Evaluate for solar mass and  $10^7$  solar mass black holes.

**Find predictions for the three classical tests of general relativity. For all three problems, assume  $\frac{2M}{r} \ll 1$ .**

1. **Perihelion advance of Mercury.** From the orbital equation derived by integrating  $\frac{dr}{d\varphi}$  (given in the Notes), find the angle between adjacent minima of  $r(\varphi)$ . The amount by which this exceeds  $2\pi$  is the perihelion advance.

2. **Gravitational red shift.** Use your expression for null geodesics, restricted to outward radial motion ( $L = 0$ ) to find the fractional change in frequency  $\frac{\omega'}{\omega}$  (or wavelength) of the light. Remember that the momentum 4-vector,  $p^\alpha = (\frac{E}{c}, \mathbf{p}) = \hbar (\frac{\omega}{c}, \mathbf{k})$  is tangent to the null curve.
  
3. **Deflection of light passing the sun.** Again consider a null geodesic but now  $L \neq 0$ . Instead,  $L$  follows from the point of closest approach,  $L = r_0 c$  where  $r_0$  is the impact parameter. The orbit is symmetric around this point, so the integral for  $r$  will run from  $\infty$  to  $r_0$  and back out again. Write the solution as  $\varphi = 2 \int_{r_0}^{\infty} F(r) dr$  where  $F(r)$  is what you get from the null geodesic equation. The deviation from a straight line is  $\Delta\varphi = 2 \int_{r_0}^{\infty} F(r) dr - \pi$ .