The geodesic equation as Newton's Law of Universal Gravitation

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1 The geodesic equation as Newton's second law

We have seen that curves of extremal proper length or time satisfy the geodesic equation,

$$\frac{d^2 x^{\alpha}}{d\tau^2} = -\Gamma^{\alpha}_{\mu\beta} u^{\mu} u^{\beta} \tag{1}$$

where $u^{\alpha} = \frac{dx^{\alpha}}{d\tau}$ is the 4-velocity. Since the acceleration may be written using the chain rule as

$$\begin{array}{rcl} \displaystyle \frac{d^2 x^{\alpha}}{d\tau^2} & = & \displaystyle \frac{du^{\alpha}}{d\tau} \\ & = & \displaystyle \frac{dx^{\beta}}{d\tau} \frac{\partial u^{\alpha}}{\partial x^{\beta}} \\ & = & \displaystyle u^{\beta} \frac{\partial u^{\alpha}}{\partial x^{\beta}} \end{array}$$

we may write the geodesic equation as

$$0 = u^{\beta} \frac{\partial u^{\alpha}}{\partial x^{\beta}} + \Gamma^{\alpha}_{\mu\beta} u^{\mu} u^{\beta}$$
$$= u^{\beta} \left(\frac{\partial u^{\alpha}}{\partial x^{\beta}} + u^{\mu} \Gamma^{\alpha}_{\mu\beta} \right)$$
$$= u^{\beta} D_{\beta} u^{\alpha}$$

This describes an *autoparallel*: the parallel transport of the vector u^{α} along its own direction.

In the form of eq.(1), the geodesic equation gives the acceleration of a particle in terms of geometric quantities. This suggests that it might be possible to explain some force using the connection. If we multiply the geodesic equation by the mass of a particle, the left side becomes the proper time rate of change of the momentum,

$$\frac{dp^{\alpha}}{d\tau} = -m\Gamma^{\alpha}_{\ \mu\beta}u^{\mu}u^{\beta}$$

where $p^{\alpha} = mu^{\alpha}$. The question is, can we make the connection term look like a force?

2 Nearly flat space

Suppose the metric is nearly that of a flat space, differing only in the time component,

$$= \left(\begin{array}{cccc} -c^2 + \phi & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{array}\right)$$

and suppose the function ϕ depends only on $r = \sqrt{x^2 + y^2 + z^2}$. Then the interval of proper time is

$$d\tau^{2} = \left(1 - \frac{1}{c^{2}}\phi\right)dt^{2} - \frac{1}{c^{2}}\left(dx^{2} + dy^{2} + dz^{2}\right)$$

We would like to find the connection, which is built from derivatives of $g_{\mu\nu}$. The only nonvanishing derivatives of the metric are

where i = 1, 2, 3 and

$$\phi_{,i} = \frac{d\phi}{dr} \frac{\partial r}{\partial x^i}$$
$$= \frac{d\phi}{dr} \frac{x_i}{r}$$

This means the only nonvanishing connection components must have two time indices, 0, and one spatial one, i:

$$\Gamma_{00i} = \Gamma_{0i0} = \frac{1}{2} (g_{00,i} + g_{0i,0} - g_{0i,0})$$
$$= \frac{1}{2} \frac{d\phi}{dr} \frac{x_i}{r}$$
$$\Gamma_{i00} = \frac{1}{2} (g_{i0,0} + g_{i0,0} - g_{00,i})$$
$$= -\frac{1}{2} \frac{d\phi}{dr} \frac{x_i}{r}$$

Raising the first index on each, we have

and

$$\begin{split} \Gamma^i_{\ 00} &= g^{i\alpha}\Gamma_{\alpha00} \\ &= g^{ij}\Gamma_{j00} \\ &= \frac{1}{2}\frac{d\phi}{dr}\frac{x^i}{r} \end{split}$$

Now return to the geodesic equation. Writing the 4-velocity as $u^{\alpha} = \gamma(c, v^i)$ and computing each component of the equation separately, we have

$$\frac{dp^0}{d\tau} = -m\Gamma^0_{\mu\beta}u^\mu u^\beta$$

$$= -m\Gamma_{00}^{0}u^{0}u^{0} - m\Gamma_{0i}^{0}u^{0}u^{i} - m\Gamma_{i0}^{0}u^{i}u^{0} - m\Gamma_{ij}^{0}u^{i}u^{j}$$

$$= -2m\Gamma_{0i}^{0}u^{0}u^{i}$$

$$= -2m\left(-\frac{1}{2(1-\phi)}\frac{d\phi}{dr}\frac{x_{i}}{r}\right)\gamma^{2}cv^{i}$$

$$= \frac{mc}{r(1-\phi)}\gamma^{2}\frac{d\phi}{dr}v^{i}x_{i}$$

for the time component and

$$\begin{aligned} \frac{dp^i}{d\tau} &= -m\Gamma^i_{\ \mu\beta}u^\mu u^\beta \\ &= -m\Gamma^i_{\ 00}u^0u^0 - m\Gamma^i_{\ 0j}u^0u^j - m\Gamma^i_{\ j0}u^ju^0 - m\Gamma^i_{\ jk}u^ju^k \\ &= -m\Gamma^i_{\ 00}u^0u^0 \\ &= -\frac{1}{2}m\gamma^2c^2\frac{d\phi}{dr}\frac{x^i}{r} \end{aligned}$$

We consider the case of non-relativistic velocities, so that we may set $\gamma \approx 1$. Then, noting that $\frac{x^2}{r}$ is just a unit vector in the radial direction, we can choose ϕ so that this last expression looks like the Newton's gravitational force,

$$-\frac{1}{2}mc^{2}\frac{d\phi}{dr}\hat{\mathbf{r}} = -\frac{GMm}{r^{2}}\hat{\mathbf{r}}$$
$$\frac{d\phi}{dr} = \frac{2GM}{r^{2}c^{2}}$$
$$\phi = -\frac{2GM}{r^{2}}$$

Integrating,

$$\phi = -\frac{2GM}{rc^2}$$

Notice that $\frac{2GM}{r}$ is just the square of the classical escape velocity,

$$\frac{1}{2}mv^2 = \frac{GMm}{r}$$
$$v^2 = \frac{2GM}{r}$$

so the magnitude of ϕ is very small, $\frac{v_{escape}^2}{c^2} \ll 1$ for planets and ordinary stars. With this choice for ϕ , the metric differs from flat space by only a term of order $\frac{v_{escape}^2}{c^2}$, and we reproduce Newton's law of gravity,

$$\frac{d\mathbf{p}}{dt} = -\frac{GMm}{r^2}\hat{\mathbf{r}}$$

at normal velocities $v \ll c$.

Now consider the time component of the geodesic equation. Substituting this result for ϕ into the 0 component of the geodesic equation,

$$\frac{dp^0}{d\tau} = \frac{1}{1 + \frac{2GM}{rc^2}} \frac{2GMm}{r^2c} \mathbf{v} \cdot \hat{\mathbf{r}}$$

Since $p^0 = \frac{E}{c}$ and $\tau \approx t$,

$$\frac{dE}{dt} = \frac{1}{1 + \frac{2GM}{rc^2}} \frac{2GMm}{r^2} \mathbf{v} \cdot \hat{\mathbf{r}}$$

$$= \frac{1}{1 + \frac{v_{escape}^2}{c^2}} \frac{2GMm}{r^2} \mathbf{v} \cdot \hat{\mathbf{r}}$$
$$\approx \frac{2GMm}{r^2} \hat{\mathbf{r}} \cdot \mathbf{v}$$
$$= 2\mathbf{F}_{Newton} \cdot \mathbf{v}$$

This differs from the actual rate of change of energy of a falling particle by a factor of 2. This factor is eliminated when we use a more realistic metric.

Exercise: Repeat these arguments using the metric

$$g_{\mu\nu} = \begin{pmatrix} -c^2 + \phi & 0 & 0 & 0 \\ 0 & 1 + \lambda \phi & 0 & 0 \\ 0 & 0 & 1 + \lambda \phi & 0 \\ 0 & 0 & 0 & 1 + \lambda \phi \end{pmatrix}$$

where λ is constant. Can you choose λ to get the correct expression for the change in energy?