

The geodesic equations are:

$$\begin{aligned}
0 &= \frac{du^0}{d\tau} - \frac{a}{r^2 \left(1 + \frac{a}{r}\right)} u^0 u^1 \\
0 &= \frac{1}{c^2} \frac{du^1}{d\tau} - \left(1 + \frac{a}{r}\right) \frac{a}{2r^2} u^0 u^0 + \frac{a}{2r^2 c^2 \left(1 + \frac{a}{r}\right)} u^1 u^1 - \frac{r}{c^2} \left(1 + \frac{a}{r}\right) u^2 u^2 - \left(1 + \frac{a}{r}\right) \frac{r}{c^2} \sin^2 \theta u^3 u^3 \\
0 &= \frac{du^2}{d\tau} - \sin \theta \cos \theta u^3 u^3 + \frac{2}{r} u^2 u^1 \\
0 &= \frac{du^3}{d\tau} + \frac{2 \cos \theta}{\sin \theta} u^2 u^3 + \frac{2}{r} u^3 u^1
\end{aligned}$$

We also have the relation given by the line element,

$$\frac{1}{c^2} \frac{ds^2}{d\tau^2} = - \left(1 + \frac{a}{r}\right) \left(\frac{dt}{d\tau}\right)^2 + \frac{1}{c^2 \left(1 + \frac{a}{r}\right)} \left(\frac{dr}{d\tau}\right)^2 + \frac{r^2}{c^2} \left(\frac{d\theta}{d\tau}\right)^2 + \frac{r^2}{c^2} \sin^2 \theta \left(\frac{d\varphi}{d\tau}\right)^2$$

Consider the case of a bound null orbit. Restrict the motion to constant r and $\theta = \frac{\pi}{2}$. Then

$$\begin{aligned}
0 &= \frac{du^0}{d\tau} \\
0 &= + \left(1 - \frac{2m}{r}\right) \frac{m}{r^2} u^0 u^0 - \left(1 - \frac{2m}{r}\right) \frac{r}{c^2} u^3 u^3 \\
0 &= \frac{du^3}{d\tau}
\end{aligned}$$

The first and third show that the 4-velocity is constant, while the second becomes

$$\begin{aligned}
0 &= \frac{mc^2}{r^3} (u^0)^2 - (u^3)^2 \\
u^3 &= \sqrt{\frac{mc^2}{r^3}} u^0
\end{aligned}$$

We also have the line element,

$$\begin{aligned}
0 &= - \left(1 - \frac{2m}{r}\right) (u^0)^2 + \frac{r^2}{c^2} (u^3)^2 \\
&= - \left(1 - \frac{2m}{r}\right) (u^0)^2 + \frac{r^2}{c^2} \frac{mc^2}{r^3} (u^0)^2 \\
0 &= -1 + \frac{2m}{r} + \frac{m}{r} \\
r &= 3m
\end{aligned}$$